

SFB F68 Tomography Across the Scales





Generalized Fourier Diffraction Theorem and Filtered Backpropagation for Tomographic Reconstruction

Michael Quellmalz | TU Berlin | GIP Meeting, Siegen, 25 September 2024 joint work with Clemens Kirisits, Eric Setterqvist





Outline

1 Introduction

2 Helmholtz Equation and Fourier Diffraction Theorem

③ Motion of the Object

4 Reconstruction of the Object





Optical Diffraction Tomography (ODT)







Optical Diffraction

Optical diffraction occurs when the wavelength of the incident wave is large \approx the size of the object (μm scale)



Simulation of the scattered field from spherical particles (size \approx wavelength)



Image with diffraction © Medizinische Universität Innsbruck



Model of Optical Diffraction Tomography (for one direction)

- We have: field $u^{\text{tot}}(\tilde{\boldsymbol{x}}, r_{\text{M}}), \tilde{\boldsymbol{x}} \in \mathbb{R}^{d-1}$, at measurement plane $x_d = r_{\text{M}}$
- We want: scattering potential f on \mathbb{R}^d with compact support
- Illumination by plane wave $u^{inc}(\mathbf{x}) = e^{ik_0\mathbf{x}\cdot\mathbf{s}}$ with direction $\mathbf{s} \in \mathbb{S}^{d-1}$ and wave number k_0
- Total field $u^{\text{tot}}(\mathbf{x}) = u(\mathbf{x}) + u^{\text{inc}}(\mathbf{x})$ solves the wave equation

$$-\left(\Delta + f(\boldsymbol{x}) + k_0^2\right) u^{\text{tot}}(\boldsymbol{x}) = 0$$

Rearranging yields

$$-\left(\Delta+k_0^2\right)u(\mathbf{x})-\underbrace{\left(\Delta+k_0^2\right)u^{\text{inc}}(\mathbf{x})}_{=0}=f(\mathbf{x})\left(u(\mathbf{x})+u^{\text{inc}}(\mathbf{x})\right)$$

Born approximation

Assuming $|u| \ll |u^{inc}|$, we obtain the Helmholtz equation

$$-\left(\Delta+k_0^2\right)u(\boldsymbol{x})=f(\boldsymbol{x})u^{\rm inc}(\boldsymbol{x})$$



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(1)

• Helmholtz equation

$$-(\Delta + k_0^2)u(\mathbf{x}) = g(\mathbf{x}), \quad \mathbf{x} \in \mathbb{R}^d$$

• The outgoing solution satisfies the Sommerfeld radiation condition

$$\lim_{r\to\infty}r^{\frac{d-1}{2}}\left(\frac{\partial u(\boldsymbol{x})}{\partial r}-\mathsf{i}k_0u(\boldsymbol{x})\right)=0$$

uniformly in **s**, where $\mathbf{x} = r\mathbf{s}$, $r = |\mathbf{x}|$, and $\partial/\partial r$ denotes the radial derivative

• It has the fundamental solution

$$G(\boldsymbol{x}) = \frac{\mathsf{i}}{4} \left(\frac{k_0}{2\pi |\boldsymbol{x}|} \right)^{\frac{d-2}{2}} H_{\frac{d-2}{2}}^{(1)}(k_0 |\boldsymbol{x}|),$$

where $H_a^{(1)}$ is the Hankel function of the first kind and order *a*.

Lemma

For $g \in L^1(\mathbb{R}^d)$ with compact support,

$$u = g * G \in L^1_{loc}(\mathbb{R}^d) \cap \mathcal{S}'(\mathbb{R}^d)$$

is the unique outgoing solution of the Helmholtz equation (1).





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Some Auxiliary Lemmas

Lemma

If $g_n \to 0$ in $L^1(\mathbb{R}^d)$ and $\bigcup_n \operatorname{supp} g_n$ is bounded, then $g_n * G \to 0$ in $\mathcal{S}'(\mathbb{R}^d)$.

- Set $\tilde{\mathcal{F}}$ as partial Fourier transform in the first d-1 coordinates
- For $\boldsymbol{x} \in \mathbb{R}^d$, set $\tilde{\boldsymbol{x}} = (x_1, \dots, x_{d-1})$
- $\kappa(\mathbf{x}) \coloneqq \sqrt{k_0^2 |\mathbf{x}|^2}$, $\mathbf{x} \in \mathbb{R}^{d-1}$

Lemma

Let $d \geq 2$. Then $\tilde{\mathcal{F}}G$ is given by the locally integrable function

$$\tilde{\mathcal{F}}G(\mathbf{x}) = (2\pi)^{\frac{1-d}{2}} \frac{\mathsf{i} \mathsf{e}^{\mathsf{i}\kappa(\tilde{\mathbf{x}})|\mathbf{x}_d|}}{2\kappa(\tilde{\mathbf{x}})}.$$

In contrast: $\mathcal{F}G$ exists as tempered distribution, but is not regular





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In contrast: $\mathcal{F}G$ exists as tempered distribution, but is not regular





Define
$$\boldsymbol{h}^{\pm}(\boldsymbol{x}) \coloneqq \begin{pmatrix} \boldsymbol{x} \\ \pm \kappa(\boldsymbol{x}) \end{pmatrix} = \begin{pmatrix} \boldsymbol{x} \\ \pm \sqrt{k_0^2 - |\boldsymbol{x}|^2} \end{pmatrix}$$
, $\boldsymbol{x} \in \mathbb{R}^{d-1}$ (hemisphere)

Generalized Fourier Diffraction Theorem

[Kirisits, Q, Setterqvist 2024]

Let $d \ge 2$ and $g \in L^1(\mathbb{R}^d)$ has compact support. Then $\tilde{\mathcal{F}}u$, where u = g * G, is given by the locally integrable function

$$\tilde{\mathcal{F}}u(\boldsymbol{x}) = \sqrt{\frac{\pi}{2}} \frac{\mathrm{i}}{\kappa(\tilde{\boldsymbol{x}})} \left(\mathrm{e}^{\mathrm{i}\kappa(\tilde{\boldsymbol{x}})_{r_d}} \mathcal{F}\left((1-\chi_{x_d})g\right)(\boldsymbol{h}^+(\tilde{\boldsymbol{x}})) + \mathrm{e}^{-\mathrm{i}\kappa(\tilde{\boldsymbol{x}})x_d} \mathcal{F}\left(\chi_{r_d}g\right)(\boldsymbol{h}^-(\tilde{\boldsymbol{x}})) \right),$$

where $\tilde{\mathbf{x}} = (x_1, ..., x_{d-1})$, $\tilde{\mathcal{F}}$ is the d-1 dimensional partial Fourier transform and χ_{x_d} is the indicator function of $\{\mathbf{y} \in \mathbb{R}^d : \mathbf{y}_d \ge x_d\}$.

If x_d is sufficiently large or sufficiently small such that

$$\pm (x_d - y_d) > 0$$
 for all $\mathbf{y} \in \operatorname{supp} g$,

then

$$\tilde{\mathcal{F}}u(\mathbf{x}) = \sqrt{\frac{\pi}{2}} \frac{\mathrm{i} \mathrm{e}^{\pm \mathrm{i}\kappa(\tilde{\mathbf{x}})x_d}}{\kappa(\tilde{\mathbf{x}})} \hat{g}(\mathbf{h}^{\pm}(\tilde{\mathbf{x}})).$$





(Classical) Fourier Diffraction Theorem

Additional assumptions:

- the incident field is a plane wave $u^{inc}(\mathbf{x}) = e^{ik_0 \mathbf{x} \cdot \mathbf{s}}$
- the measurement plane $x_d = r_M$ not intersect supp f

Then

$$\sqrt{\frac{2}{\pi}}\kappa i e^{-i\kappa m} \tilde{\mathcal{F}} \underbrace{u(\tilde{\boldsymbol{x}}, r_{m})}_{\text{measured}} = \mathcal{F}f(\boldsymbol{h}(\tilde{\boldsymbol{x}}) - k_{0}\boldsymbol{s}), \quad \tilde{\boldsymbol{x}} \in \mathbb{R}^{d-1},$$

where $\boldsymbol{h}(\tilde{\boldsymbol{x}}) \coloneqq \begin{pmatrix} \tilde{\boldsymbol{x}} \\ \kappa \end{pmatrix}$ and $\kappa = \sqrt{k_0^2 - |\tilde{\boldsymbol{x}}|^2}$.

Formula well-known from [Wolf 1969] [Natterer Wuebbeling 2001] [Kak Slaney 2001]



Semisphere h(k) of available data in Fourier space





Focused Beams

- The generalized version allows for other incidence fields
- For an incident Herglotz wave

$$u^{\rm inc}(\textbf{\textit{r}}) = \int_{\mathbb{S}^{d-1}} a(\textbf{\textit{s}}) {\rm e}^{{\rm i} k_0 \textbf{\textit{s}} \cdot \textbf{\textit{r}}} \, {\rm d} \textbf{\textit{s}}(\textbf{\textit{s}}), \quad {\rm where} \ \textbf{\textit{a}} \in L^2(\mathbb{S}^{d-1}),$$

it yields

$$\tilde{\mathcal{F}}u(\mathbf{x},\mathbf{r}_{\mathrm{M}}) = \sqrt{\frac{\pi}{2}} \frac{\mathrm{i}\mathrm{e}^{\pm\mathrm{i}\kappa(\mathbf{x})\mathbf{r}_{\mathrm{M}}}k_{0}^{2}}{\kappa(\mathbf{x})} \int_{\mathbb{S}^{d-1}} a(\mathbf{s})\hat{f}(\mathbf{h}^{\pm}(\mathbf{x}) - k_{0}\mathbf{s})\,\mathrm{d}s(\mathbf{s}).$$

• Reconstruction algorithms for focused beams (see images below) in [Kirisits Naujoks Scherzer 2024]







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Moving the Object and Incidence Direction

- Rigid motion of the object causes scattering potential $f(R_t(\mathbf{x} \mathbf{d}_t))$ with
 - Rotation $R_t \in SO(d)$ (with $R_0 := id$)
 - Translation $\boldsymbol{d}_t \in \mathbb{R}^d$ (with $\boldsymbol{d}_0 \coloneqq \mathbf{0}$)
- Incidence direction $\boldsymbol{s}_t \in \mathbb{S}^{d-1}$
- Position of the measurement plane is kept fixed (equivalent to moving object and incidence simultaneously)

Fourier diffraction theorem (with motion)

The quantity

$$\mu_t(\mathbf{x}) \coloneqq \sqrt{\frac{2}{\pi}} \kappa \mathrm{i} \mathrm{e}^{-\mathrm{i}\kappa r_{\mathrm{M}}} \tilde{\mathcal{F}} u(\mathbf{k}, r_{\mathrm{M}}) = \mathcal{F} f(\underbrace{\mathbf{R}_t \mathbf{h}(\mathbf{x}) - k_0 \mathbf{s}_t}_{\text{Fourier cover}}) \frac{\mathrm{e}^{-\mathrm{i}\langle \mathbf{d}_t, \mathbf{h}(\mathbf{x}) \rangle}}{\mathrm{e}^{-\mathrm{i}\langle \mathbf{d}_t, \mathbf{h}(\mathbf{x}) \rangle}}, \quad \|\mathbf{x}\| < k_0,$$

depends only on the measurements of u.



Fourier Cover: Angle Scan



2D Fourier coverage for incidence direction $\mathbf{s}(t) = (\cos t, \sin t)$. Measurements are taken at $r_2 = r_M$. The Fourier coverage (light red) is a union of infinitely many semicircles, whose centers lie on the dashed blue curve.







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Figure: 2D Fourier coverage for a rotating object, incidence direction $\mathbf{s} = (0, 1)$ and measurements taken at $r_2 = r_M$. The Fourier coverage (light red) is a union of infinitely many semicircles, some of which are depicted in red. (d): 3D Fourier coverage for incidence $\mathbf{s} = (0, 0, 1)$ and rotation around first axis.



Fourier Cover: Object Rotation (Incidence Perpendicular to Measurement Plane)



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Figure: 2D Fourier coverage for a rotating object, incidence direction s = (1,0) and measurements taken at $r_2 = r_M$.



Figure: 3D Fourier coverage for a full rotation of the object about the r_1 -axis with incidence direction $\mathbf{s} = (0, 1, 0)$.





Fourier cover: Angle scan & Rotation





3D Fourier cover

- Incidence is rotated along a half circle
- and the experiment is repeated with the object rotated by $90^\circ.$





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Apprach 1: Filtered Backpropagation

Idea: Inverse Fourier transform of $\mathcal{F}f$ restricted to the set of available data \mathcal{Y} ,

$$\mathbf{f}_{\mathrm{bp}}(\mathbf{r}) := (2\pi)^{-\frac{d}{2}} \int_{\mathcal{Y}} \mathcal{F}f(\mathbf{y}) \, \mathrm{e}^{\mathrm{i}\mathbf{y}\cdot\mathbf{r}} \, \mathrm{d}\mathbf{y}$$

with the transformation $T(\mathbf{x}, t) \coloneqq R_t \mathbf{h}(\mathbf{x})$

Theorem

Let the rotation $R_t \in SO(d)$, translation d_t and incidence $s_t \in \mathbb{S}^{d-1}$ be piecewice C^1 . Then

$$f_{\mathsf{bp}}(\mathbf{r}) = (2\pi)^{-\frac{d}{2}} \int_0^T \int_{\mathcal{B}_{k_0}} \mathcal{F}f(T(\mathbf{x}, t)) \, \mathsf{e}^{\mathsf{i} \, T(\mathbf{x}, t) \cdot (\mathbf{r} + d_t)} \, \frac{|\det \nabla T(\mathbf{x}, t)|}{\operatorname{Card} \, T^{-1}(T(\mathbf{x}, t))} \, \mathsf{d}\mathbf{x} \, \mathsf{d}t,$$
$$\nabla T(\mathbf{x}, t) = \frac{k_0(t)k_0'(t) - R_t \mathbf{h}(\mathbf{x}, t) \, (k_0(t)R_t \mathbf{s}_t)'}{\kappa}.$$

Banach indicatrix $Card(T^{-1}(\mathbf{y}))$ needs to be estimated (except for special cases). Well-known for rotation around coordinate axis [Devaney 1982]







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$$\mathrm{e}\, \det \nabla\mathcal{T}(\mathbf{x},t) = \frac{k_0(t)k_0'(t) - R_t \mathbf{h}(\mathbf{x},t) \, (k_0(t)R_t \mathbf{s}_t)'}{\kappa}.$$

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[Kirisits, Q, Setterqvist 2024]









Discretization

- Object $f(\mathbf{x}_{\mathbf{k}})$ with $\mathbf{x}_{\mathbf{k}} = \mathbf{k} \frac{2L_s}{\kappa}$, $\mathbf{k} \in \mathcal{I}_{\kappa}^d \coloneqq \{-\kappa/2, \dots, \kappa/2 1\}^d$
- Measurements $u_{t_m}^{\text{tot}}(\boldsymbol{y}_n, r_M)$ with $\boldsymbol{y}_n = \boldsymbol{n} \frac{2L_M}{N}$, $\boldsymbol{n} \in \mathcal{I}_N^{d-1}$
- discrete Fourier transform (DFT)

$$[\boldsymbol{F}_{\mathsf{DFT}} u_{t_m}]_{\boldsymbol{\ell}} := \sum_{\boldsymbol{n} \in \mathcal{I}_N^{d-1}} u_{t_m}(\boldsymbol{y}_{\boldsymbol{n}}, r_{\mathsf{M}}) \, \mathrm{e}^{-2\pi \mathrm{i} \boldsymbol{n} \cdot \boldsymbol{\ell} / N}, \qquad \boldsymbol{\ell} \in \mathcal{I}_N^{d-1},$$

• Non-uniform discrete Fourier transform (NDFT)

$$[\mathbf{F}_{\mathsf{NDFT}}\mathbf{f}]_{m,\boldsymbol{\ell}} \coloneqq \sum_{\mathbf{k}\in\mathcal{I}_{K}^{d}} f_{\mathbf{k}} \, \mathrm{e}^{-\mathrm{i}\mathbf{x}_{\mathbf{k}}\cdot\left(R_{i_{m}}\mathbf{h}(\mathbf{y}_{\boldsymbol{\ell}})\right)}, \qquad m\in\mathcal{J}_{M}, \ \boldsymbol{\ell}\in\mathcal{I}_{N}^{d-1}$$

Discretized forward operator

$$oldsymbol{D}^{ ext{tot}}oldsymbol{f}\coloneqqoldsymbol{F}_{ ext{DFT}}(oldsymbol{c}\odotoldsymbol{F}_{ ext{NDFT}}oldsymbol{f})+ extbf{e}^{ ext{i}k_0r_{ ext{M}}}, \qquadoldsymbol{f}\in\mathbb{R}^{K^d},$$

where $\boldsymbol{c} = \left[\frac{i}{\kappa(\boldsymbol{y}_{\ell})} e^{i \kappa(\boldsymbol{y}_{\ell}) \boldsymbol{f}_{M}} \left(\frac{N}{L_{M}}\right)^{d-1} \left(\frac{L_{s}}{K}\right)^{d}\right]_{\ell \in \mathcal{I}_{M}^{d-1}}$



1



Reconstruction of f

Inverse

$$m{r} pprox m{F}_{ extsf{NDFT}}^{-1} ig((m{F}_{ extsf{DFT}}m{u}^{ extsf{tot}} - extsf{e}^{ extsf{i}k_0m{r}_{ extsf{M}}}ig) \oslash m{c} ig)$$

Crucial part: inversion of NDFT $\mathbf{F}_{\text{NDFT}}^{-1}$





Approach 2: Conjugate Gradient (CG) Method

• Conjugate Gradients (CG) on the normal equations

$$\underset{\boldsymbol{f} \in \mathbb{R}^{k^3}}{\operatorname{arg\,min}} \quad \|\boldsymbol{F}_{\mathsf{NDFT}}(\boldsymbol{f}) - \boldsymbol{g}\|_2^2$$

• NFFT (Non-uniform fast Fourier transform) for computing $F_{NDFT}(f)$ in $O(N^3 \log N)$ steps [Dutt Rokhlin 93], [Beylkin 95], [Potts Steidl Tasche 01], [Potts Kunis Keiner 04+]

Approach 3: TV (Total Variation) Regularization

• Regularized inverse

$$\underset{t \in \mathbb{R}^{K^3}}{\operatorname{arg\,min}} \qquad \chi_{\mathbb{R}^{K^3}_{\geq 0}}(f) + \tfrac{1}{2} \| \mathbf{F}_{\mathsf{NDFT}}(f) - \boldsymbol{g} \|_2^2 + \lambda \mathsf{TV}(f),$$

- Primal-dual (PD) iteration [Chambolle & Pock 2010]
- Adaptive selection of step sizes [Yokota & Hontani 2017]



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3D Reconstruction: Moving Rotation Axis and 5% Gaussian Noise



PSNR 25.50 SSIM 0.157

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Conclusions

- Generalized Fourier diffraction theorem on $L^1(\mathbb{R}^d)$
- Filtered backpropagation formula for a wide range of experimental setups
- Compared image reconstruction methods





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Thank you for your attention!





References of This Talk



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Estimation of the Banach indicatrix

- Banach indicatrix $\operatorname{Card}(\mathcal{T}_{\pm}^{-1}(\boldsymbol{y}))$ counts how often a point $\boldsymbol{y} \in \mathbb{B}^{d}_{2k_{0}}$ is "hit" by \mathcal{T}_{+} .
- In a discrete setting, it is unlikely that a point y is exactly hit by any grid point $T_+(\tilde{x}_m, t_n)$.
- The coverage for fixed time t is the hemisphere

$$\{T_{+}(\tilde{\boldsymbol{x}},t):\tilde{\boldsymbol{x}}\in\mathbb{B}_{k_{0}}^{d-1}\}=\left\{\boldsymbol{z}\in\mathbb{R}^{d}:\|\boldsymbol{z}+k_{0}(t)R_{t}\boldsymbol{s}_{t}\|=k_{0}(t),\,\boldsymbol{z}\cdot R_{t}\boldsymbol{e}^{d}>-k_{0}(t)\boldsymbol{s}_{t}\cdot\boldsymbol{e}^{d}\right\}$$

which moves continuously with t.

• Idea: a point y is hit by T_+ between the time steps t_{n-1} and t_n , if the sign of $||y - R_t s_t|| - k_0(t)$ changes between these time steps. We use the approximation

$$\operatorname{Card}(\mathcal{T}_{+}^{-1}(\boldsymbol{y})) \approx \sum_{n=1}^{N} \frac{\chi(n)}{2} \left| \operatorname{sgn}(\|\boldsymbol{y} + k_0(t_{n-1}) \boldsymbol{R}_{t_{n-1}} \boldsymbol{s}_{t_{n-1}}\| - k_0(t_{n-1})) - \operatorname{sgn}(\|\boldsymbol{y} + k_0(t_n) \boldsymbol{R}_{t_n} \boldsymbol{s}_{t_n}\| - k_0(t_n)) \right|,$$

where

$$\chi(\mathbf{n}) \coloneqq \begin{cases} 1, & \text{if } \mathbf{y} \cdot (\mathbf{R}_t \mathbf{e}^d) > -k_0(t_n) \mathbf{s}_{t_n} \cdot \mathbf{e}^d, \\ 0, & \text{otherwise.} \end{cases}$$

with the sign function $\operatorname{sgn}(x) \coloneqq \frac{x}{|x|}$ for $x \neq 0$ and $\operatorname{sgn}(0) \coloneqq 0$.