



Approximation Properties of the Double Fourier Sphere Method

Michael Quellmalz | TU Berlin | Waves Conference, Karlsruhe, 15 February 2022 joint work with Sophie Mildenberger





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Introduction

Analysis

Fourier series of DFS functions

Numerical example

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Spherical functions

Sphere $\mathbb{S}^2 = \{ \boldsymbol{\xi} \in \mathbb{R}^3 : ||\boldsymbol{\xi}|| = 1 \}$ Function $f : \mathbb{S}^2 \to \mathbb{C}$



Approximation of f

Write as finite sum of basis functions

$$f(\boldsymbol{\xi}) = \sum_{n=0}^{N} f_n b_n(\boldsymbol{\xi})$$

Choice of basis functions:

- Spherical harmonics
- Trigonometric polynomials (this talk)

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Spherical harmonics

Use spherical coordinates

 $\phi(\lambda,\theta) = (\cos\lambda\,\sin\theta,\sin\lambda\,\sin\theta,\cos\theta)$

Define the spherical harmonics of degree n

$$Y_n^k(\phi(\lambda,\theta)) = \sqrt{\frac{2n+1}{4\pi} \frac{(n-k)!}{(n+k)!}} P_n^k(\cos\theta) e^{ik\lambda}$$

 P_n^k .. associated Legendre function

- Y_n^k are orthogonal polynomials of degree n
- Y_n^k are eigenfunctions of the spherical Laplacian and Funk-Radon transforms
- fast numerical transforms available (NFSFT, SHT) [Driscoll, Healy 1994] [Kunis, Potts 2003] [Schaeffer 2013]
 However: Computations are still considerably slower than classical FFT on the torus





Fourier series on the torus

Torus $\mathbb{T}^2 := [-\pi, \pi)^2$. Function $g: \mathbb{T}^2 \to \mathbb{C} \iff g: \mathbb{R}^2 \to \mathbb{C}$ is 2π -biperiodic

Fourier series on the torus

Let $g \in L_2(\mathbb{T}^2)$ and $\mathbf{n} \in \mathbb{Z}^2$. We define the *n*-th Fourier coefficient

$$c_{\mathbf{n}}(\mathbf{g}) := (2\pi)^{-2} \int_{\mathbb{T}^2} \mathbf{g}(\mathbf{x}) \, \mathrm{e}^{-\mathrm{i} \langle \mathbf{n}, \mathbf{x} \rangle} \, \mathrm{d}\mathbf{x},$$

the N-th partial Fourier sum

$$F_N g(\mathbf{x}) := \sum_{|\mathbf{n}|_{\infty} \leq N} c_{\mathbf{n}}(g) e^{\mathrm{i} \langle \mathbf{n}, \mathbf{x} \rangle}, \qquad \mathbf{x} \in \mathbb{T}^2,$$

and the Fourier series

$$Fg := \lim_{N \to \infty} F_N g.$$

 F_{Ng} can be computed efficiently with the fast Fourier transform (FFT) in $\mathcal{O}(N \log N)$ steps

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The Double Fourier Sphere (DFS) method I

Spherical coordinates

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 $\phi(\lambda,\theta) = (\cos\lambda\,\sin\theta,\sin\lambda\,\sin\theta,\cos\theta), \qquad \lambda \in [-\pi,\pi), \theta \in [0,\pi]$



longitude-latitude



 2π periodic in λ , but not periodic in θ





The Double Fourier Sphere (DFS) method II

Spherical coordinates

 $\phi(\lambda,\theta) = (\cos\lambda\,\sin\theta,\sin\lambda\,\sin\theta,\cos\theta), \qquad \lambda \in [-\pi,\pi), \theta \in [-\pi,\pi]$







Properties of the DFS function \tilde{f}

- 1. \tilde{f} is block-mirror centrosymmetric (BMC) $\tilde{f}(\lambda, \theta) = \tilde{f}(\lambda + \pi, -\theta)$
- 2. $\tilde{f}(\cdot, k\pi)$ is constant for all $k \in \mathbb{Z}$
- 3. \tilde{t} is 2π -biperiodic (i.e. a function on \mathbb{T}^2)

Corollary

For every function $g: \mathbb{T}^2 \to \mathbb{C}$ that admints 1. and 2., there exists a unique $f: \mathbb{S}^2 \to \mathbb{C}$, such that $g = \tilde{f} = f \circ \phi$.







The DFS method and Fourier series

Expand DFS function $\tilde{t} = t \circ \phi$ as 2D Fourier series

$$\tilde{f}(\lambda, heta) = \sum_{\mathbf{n} \in \mathbb{Z}^2} c_{\mathbf{n}}(\tilde{f}) e^{\mathrm{i}(n_1 \lambda + n_2 \phi)}.$$

Our contribution

- Show that \tilde{f} has a "similar" smoothness as f and
- can be approximated by a certain Fourier series, which converges to f.

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Function spaces

Continuously differentiable functions. Let $U \subset \mathbb{R}^d$ open and $k \in \mathbb{N}_0$. We define

$$\mathcal{C}^{k}(U) := \{ f \in \mathcal{C}(U) \mid D^{\beta}f \in \mathcal{C}(U) \text{ for all } \beta \in \mathbb{N}_{0}^{d}, \, |\beta| \leq k \}.$$

with the norm

$$\|f\|_{\mathcal{C}^{k}(U)} := \sum_{|\boldsymbol{\beta}| \leq k} \sup_{\boldsymbol{x} \in U} \left| D^{\boldsymbol{\beta}} f(\boldsymbol{x}) \right|.$$

On \mathbb{S}^2 : Definition via \mathcal{C}^k extension to \mathbb{R}^3 .

Hölder space. For $f \in \mathcal{C}^k(U)$ and $\alpha \in (0, 1)$, we define the $\mathcal{C}^{k, \alpha}$ -seminorm

$$|f|_{\mathcal{C}^{k,\alpha}(U)} := \sup_{\substack{\mathbf{x},\mathbf{y}\in U, \, \mathbf{x}\neq\mathbf{y}\\ \boldsymbol{\beta}\in\mathbb{N}_0^d, \, |\boldsymbol{\beta}|=k}} \frac{\|D^{\boldsymbol{\beta}}f(\mathbf{x}) - D^{\boldsymbol{\beta}}(\mathbf{y})\|}{\|\mathbf{x}-\mathbf{y}\|^{\alpha}}.$$

The (k, α) -Hölder space

 $\mathcal{C}^{k,\alpha}(U) := \{ f \in \mathcal{C}^k(U) \mid |f|_{\mathcal{C}^{k,\alpha}(U)} < \infty \}$

is a Banach space equipped with the norm $\|\cdot\|_{\mathcal{C}^{k,\alpha}(U)} := \|\cdot\|_{\mathcal{C}^{k}(U)} + |\cdot|_{\mathcal{C}^{k,\alpha}(U)}$.





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The (k, α) -Hölder space

$$\mathcal{C}^{^{k,\alpha}}(U) := \{ f \in \mathcal{C}^{^{k}}(U) \mid |f|_{\mathcal{C}^{k,\alpha}(U)} < \infty \}$$

is a Banach space equipped with the norm $\|\cdot\|_{\mathcal{C}^{k,\alpha}(\mathcal{U})}:=\|\cdot\|_{\mathcal{C}^{k}(\mathcal{U})}+|\cdot|_{\mathcal{C}^{k,\alpha}(\mathcal{U})}.$



Transfer of Hölder smoothness

Theorem

[Mildenberger, Q. 2021]

Let $k \in \mathbb{N}_0$ and $f \in \mathcal{C}^{k+1}(\mathbb{S}^2)$. Then for all $0 < \alpha < 1$, the DFS function \tilde{f} of f is in $\mathcal{C}^{k,\alpha}(\mathbb{T}^2)$ with

$$\tilde{f}\Big|_{\mathcal{C}^{k,\alpha}(\mathbb{T}^2)} \le (k+3)! \, \|f\|_{\mathcal{C}^{k+1}(\mathbb{S}^2)},$$

and

$$\left\| \tilde{f} \right\|_{\mathcal{C}^{k,\alpha}(\mathbb{T}^2)} \leq \frac{7}{4} \left(k+3 \right)! \left\| f \right\|_{\mathcal{C}^{k+1}(\mathbb{S}^2)}.$$

Theorem (Hölder space)

[Mildenberger, Q. 2021]

Let $k \in \mathbb{N}_0$, $0 < \alpha < 1$ and $f \in \mathcal{C}^{k, \alpha}(\mathbb{S}^2)$. Then the DFS function \tilde{f} of f is in $\mathcal{C}^{k, \alpha}(\mathbb{T}^2)$ with

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$$\left|\tilde{f}\right|_{\mathcal{C}^{k,\alpha}(\mathbb{T}^2)} \leq (k+3)! \, \|f\|_{\mathcal{C}^{k,\alpha}(\mathbb{S}^2)}$$

and

$$\left\|\tilde{t}\right\|_{\mathcal{C}^{k,\alpha}(\mathbb{T}^2)} \leq \frac{7}{4} \left(k+3\right)! \|t\|_{\mathcal{C}^{k,\alpha}(\mathbb{S}^2)}.$$

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Sobolev smoothness

Does the result for Hölder spaces also hold for Sobolev spaces?

Spherical Sobolev space

$$H^{s}(\mathbb{S}^{2}) := \left\{ f \in L^{2}(\mathbb{S}^{2}) : \|f\|_{H^{s}(\mathbb{S}^{2})} = \|(-\Delta + \frac{1}{4})^{s/2}f\|_{L^{2}(\mathbb{S}^{2})} < \infty \right\}.$$

Theorem (Sobolev space) Let $\left(\ln \left(\ln - \pi \right) \right)$

$$f: \mathbb{S}^2 \to \mathbb{R}, \ f(\xi) = \begin{cases} \ln\left(\ln\frac{1}{\sqrt{1-\xi_3^2}}\right), & |\xi_3| \neq 1, \\ 0, & \text{otherwise.} \end{cases}$$

Then $f \in H^1(\mathbb{S}^2)$, but $\tilde{f} \notin H^1(\mathbb{T}^2)$.

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DFS Fourier series

- Represent $f: \mathbb{S}^2 \to \mathbb{C}$ via the Fourier series of its DFS function $\tilde{f}: \mathbb{T}^2 \to \mathbb{C}$,

$$F_N \tilde{f}(\boldsymbol{x}) = \sum_{|\boldsymbol{n}|_{\infty} \leq N} c_{\boldsymbol{n}}(\tilde{f}) e^{i\langle \boldsymbol{n}, \boldsymbol{x} \rangle}, \qquad \boldsymbol{x} \in \mathbb{T}^2.$$

- $\tilde{\textit{f}}$ is BMC-symmetric, but the basis functions $\mathrm{e}^{\mathrm{i}(n_1\lambda+n_2\theta)}$ are not BMC
- Define BMC-symmetric function

$$e_n(\lambda,\theta) := \begin{cases} e^{i(n_1\lambda+n_2\theta)} + (-1)^{n_1}e^{i(n_1\lambda-n_2\theta)}, & n_2 \neq 0, \\ e^{i(n_1\lambda)}, & n_2 = 0. \end{cases}$$

- Define the basis functions

$$b_n(\boldsymbol{\xi}) \coloneqq e_n(\phi^{-1}(\boldsymbol{\xi})), \qquad n \in \mathbb{Z} \times \mathbb{N}_0, \ \boldsymbol{\xi} \in \mathbb{S}^2.$$

 $(\phi \text{ is bijective on } [-\pi,\pi) \times (0,\pi) \cup \{(0,0),(0,\pi)\})$





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- Represent $f: \mathbb{S}^2 \to \mathbb{C}$ via the Fourier series of its DFS function $\tilde{f}: \mathbb{T}^2 \to \mathbb{C}$,

$$F_N \tilde{f}(\mathbf{x}) = \sum_{|\mathbf{n}|_{\infty} \leq N} c_{\mathbf{n}}(\tilde{f}) e^{i\langle \mathbf{n}, \mathbf{x} \rangle}, \qquad \mathbf{x} \in \mathbb{T}^2.$$

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- Define BMC-symmetric function

$$\mathbf{e}_{\mathbf{n}}(\lambda,\theta) := \begin{cases} \mathsf{e}^{\mathsf{i}(n_1\lambda+n_2\theta)} + (-1)^{n_1} \mathsf{e}^{\mathsf{i}(n_1\lambda-n_2\theta)}, & n_2 \neq 0, \\ \mathsf{e}^{\mathsf{i}(n_1\lambda)}, & n_2 = 0. \end{cases}$$

- Define the basis functions

$$m{b}_{m{n}}(m{\xi})\coloneqq m{e}_{m{n}}(\phi^{-1}(m{\xi})), \qquad m{n}\in\mathbb{Z} imes\mathbb{N}_0,\ m{\xi}\in\mathbb{S}^2.$$

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DFS Fourier series II

Set

$$\tilde{L}_2(\mathbb{S}^2) := \left\{ f \colon \mathbb{S}^2 \to \mathbb{C} : \|f\|_{\tilde{L}_2(\mathbb{S}^2)}^2 := \int_{\mathbb{S}^2} |f(\boldsymbol{\xi})|^2 \frac{1}{\sqrt{1-\xi_3^2}} \, \mathrm{d}\boldsymbol{\xi} < \infty \right\}.$$

Theorem

[Mildenberger, Q. 2021]

Let $f \in \tilde{L}_2(\mathbb{S}^2)$. Then the Fourier coefficients of its DFS function $\tilde{f} \in L_2(\mathbb{T}^2)$ satisfy

$$\boldsymbol{c}_{n_1,n_2}(\tilde{f}) = (-1)^{n_1} \boldsymbol{c}_{n_1,-n_2}(\tilde{f}), \qquad \boldsymbol{n} \in \mathbb{Z}^2.$$

Then we have

$$S_N f(\boldsymbol{\xi}) := \sum_{n_1=-N}^N \sum_{n_2=0}^N c_n(\tilde{f}) \, b_n(\boldsymbol{\xi}) = F_N \tilde{f}(\phi^{-1}(\boldsymbol{\xi})), \qquad \boldsymbol{\xi} \in \mathbb{S}^2.$$

The set $\{b_n \mid n \in \mathbb{Z} \times \mathbb{N}_0\}$ is an orthogonal basis of $\tilde{L}_2(\mathbb{S}^2)$.





Convergence of DFS Fourier series

Known on the torus \mathbb{T}^2

If $g \in \mathcal{C}^{k,\alpha}(\mathbb{T}^2)$, its Fourier series converges with $\|f - F_N f\|_{\mathcal{C}(\mathbb{T}^2)} \le M_{k,\alpha} \|f\|_{\mathcal{C}^{k,\alpha}(\mathbb{T}^2)} N^{1-k-\alpha}$.

Theorem

[Mildenberger, Q. 2021]

Let $k \in \mathbb{N}$, $0 < \alpha \leq 1$, and $f \in C^{k,\alpha}(\mathbb{S}^2)$. Then there is a constant $M_{k,\alpha}$ depending only on k and α such that

$$\left\|f-\sum_{n_1=-N}^N\sum_{n_2=0}^N c_n(\tilde{f})b_n\right\|_{\mathcal{C}(\mathbb{S}^2)} \le M_{k,\alpha} \|f\|_{\mathcal{C}^{k,\alpha}(\mathbb{S}^2)} N^{1-k-\alpha}.$$





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DFS Fourier expansion

spherical harmonics expansion

Numerical example

Compare DFS Fourier expansion with spherical harmonics expansion



Logarithmic plot of the maximal error $\|f - S_N f\|_{\mathcal{C}(\mathbb{S}^2)}$ Computation time for expansion of degree N = 1024:

- DFS Fourier expansion 0.06 s
- spherical harmonic expansion $0.67 \,\mathrm{s}$ with NFFT (+ precomputations $3.05 \,\mathrm{s}$)

Fourier expansion $S_N f$ consists of (N + 1)(2N + 1) summands, while the spherical harmonics expansion consists of $(N + 1)^2$ summands.



DFS function $\tilde{f} \in \mathcal{C}^{3,\alpha}(\mathbb{T}^2)$





Comparison

	DFS Fourier expansion	Spherical harmonics expansion
Basis functions	$ \begin{aligned} \mathbf{e}^{\mathbf{i}(n_1\lambda+n_2\theta)} + (-1)^{n_1} \mathbf{e}^{\mathbf{i}(n_1\lambda-n_2\theta)}, \\ n_1 \in \mathbb{Z}, \ n_2 \in \mathbb{N}_0 \end{aligned} $	$Y_n^k = \mathcal{P}_n^k(\cos\theta) \mathrm{e}^{\mathrm{i}k\lambda}, \ n \in \mathbb{N}_0, \ k = -n,, n$
Orthogonality	$ ilde{\mathcal{L}}^2(\mathbb{S}^2)$ with weight $\textit{w}(\lambda, heta) = rac{1}{\sin heta}$	$L^2(\mathbb{S}^2)$
Rotational invariance	×	✓ (for fixed degree n)
Eigenfunctions	×	of the Laplace–Beltrami operator Δ
Computation	via 2D-FFT	via fast Legendre transform and (no tensor product structure, slower than 2D-FFT)
Convergence	1	✓







- DFS method maps the Hölder space $\mathcal{C}^{k,\alpha}(\mathbb{S}^2)$ continuously into the respective spaces on the torus \mathbb{T}^2 .
- We provide a series expansion of the spherical function in terms of basis functions that are orthogonal with respect to a weight on the sphere.
- The DFS Fourier series of Hölder functions converges with the expected rate.
- Numerical tests indicate a comparable approximation quality as a spherical harmonics expansion.
- DFS method is faster than spherical harmonics expansion.
- DFS method appends to other geometries.

Thank you for your attention!







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