



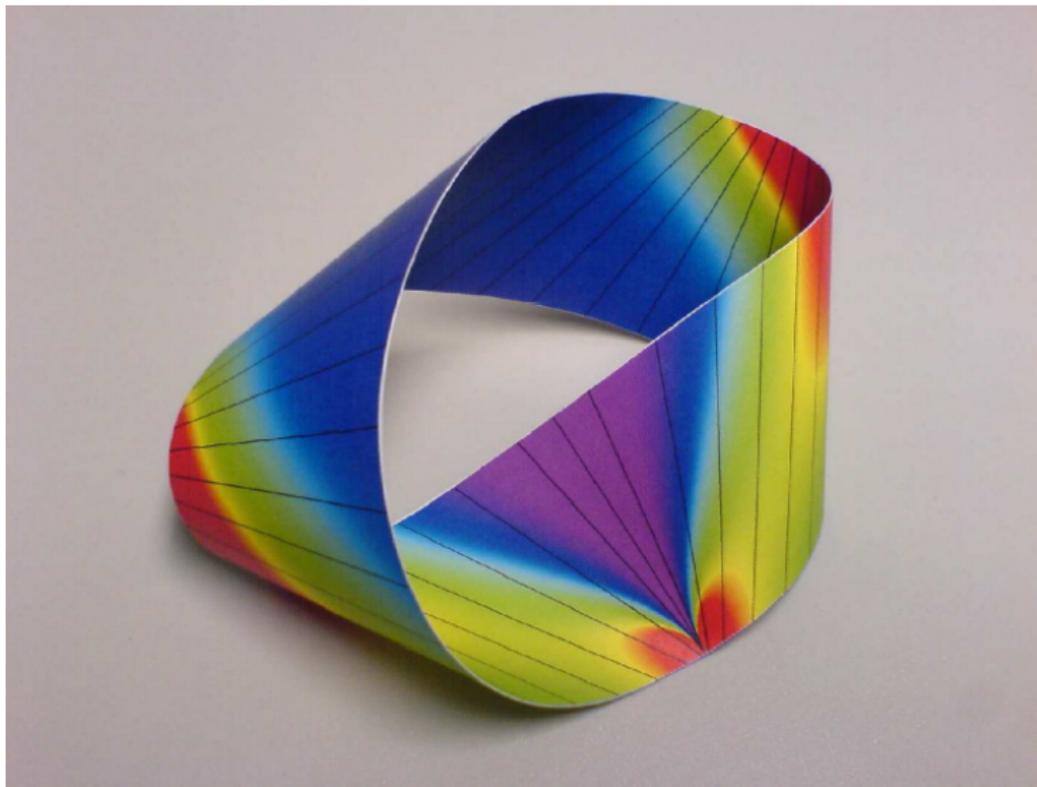
Elastic strips

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Mathematics for key technologies



May 1, 2010





- ▶ Let $\gamma : [0, L] \rightarrow \mathbb{R}^3$ be an arclength parametrized curve with Frenet frame T, N, B . Assume that there is a smooth function $\lambda : [0, L] \rightarrow \mathbb{R}$ such that the curvature and torsion of γ satisfy

$$\tau = \lambda\kappa.$$

- ▶ Then $f : [0, L] \times [-\epsilon, \epsilon] \rightarrow \mathbb{R}^3$ defined by

$$f(s, t) = \gamma(s) + t(B(s) + \lambda T(s))$$

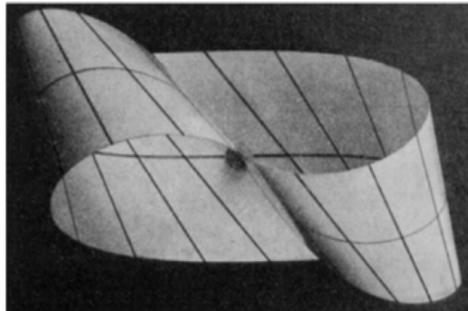
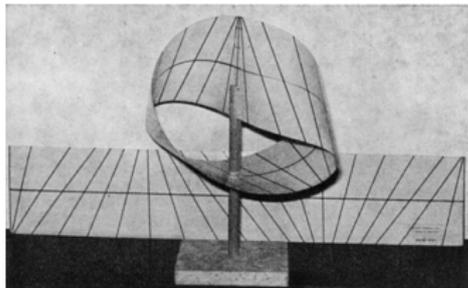
parametrizes a developpable strip of width 2ϵ .



- ▷ α the angle between the curve and the normal to the rulings \rightsquigarrow

$$\lambda = \tan \alpha$$

- ▷ Photographs 1961 by W. Wunderlich





- ▶ The bending energy

$$E = \int H^2 dA$$

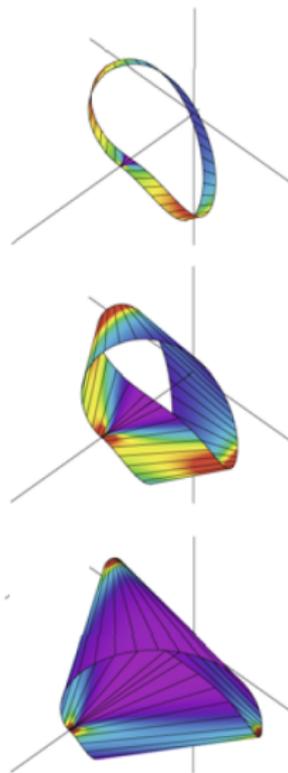
can be expressed as

$$E = \int_0^L \kappa^2 (1 + \lambda^2)^2 \frac{\log(1 + \epsilon\lambda') - \log(1 - \epsilon\lambda')}{\lambda'} ds.$$

- ▶ Wunderlich 1961

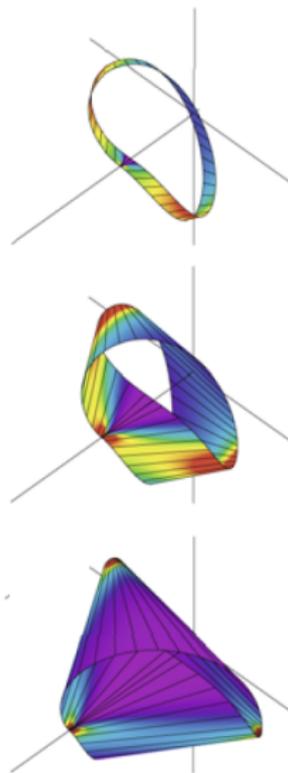


- ▶ Starostin and van der Heijden computed in 2007 the Euler-Lagrange equations
- ▶ Variational complex \rightsquigarrow Computer algebra
- ▶ Numerical study of energy-minimizing Moebius bands



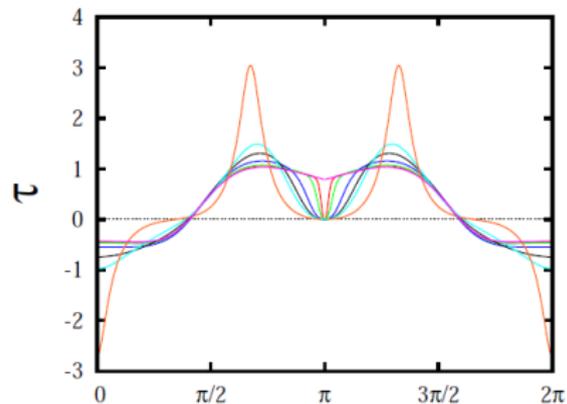
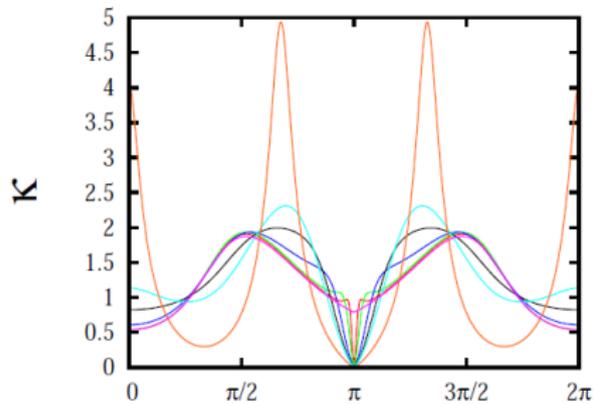


- ▶ In the limit $\epsilon \rightarrow 0$ the curvature κ jumps from $+1$ to -1 at the inflection point.
- ▶ Torsion τ is continuous with value 1.
- ▶ Angle between curve and rulings jumps from 45° to -45° .
- ▶ Energy minimizing Moebius bands are only C^1 .



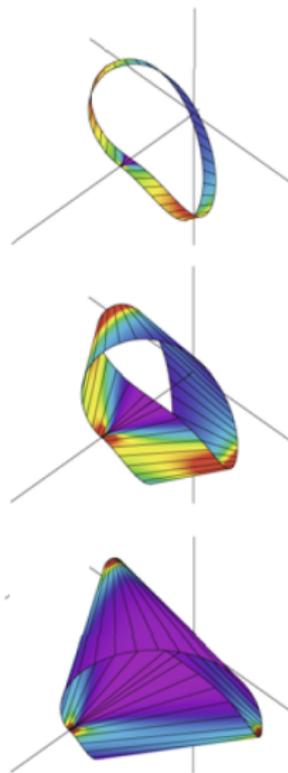


Infinitesimally thin bands





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- ▷ In the limit $\epsilon \rightarrow 0$ the energy becomes

$$E = \int_0^L \kappa^2 (1 + \lambda^2)^2 ds.$$

- ▷ Introduced by Sadowski in 1930 as

$$E = \int_0^L \frac{(\kappa^2 + \tau^2)^2}{\kappa^2} ds.$$

- ▷ E has to be minimized among strips with fixed length \rightsquigarrow Lagrange multiplier μ in Euler-Lagrange equations.

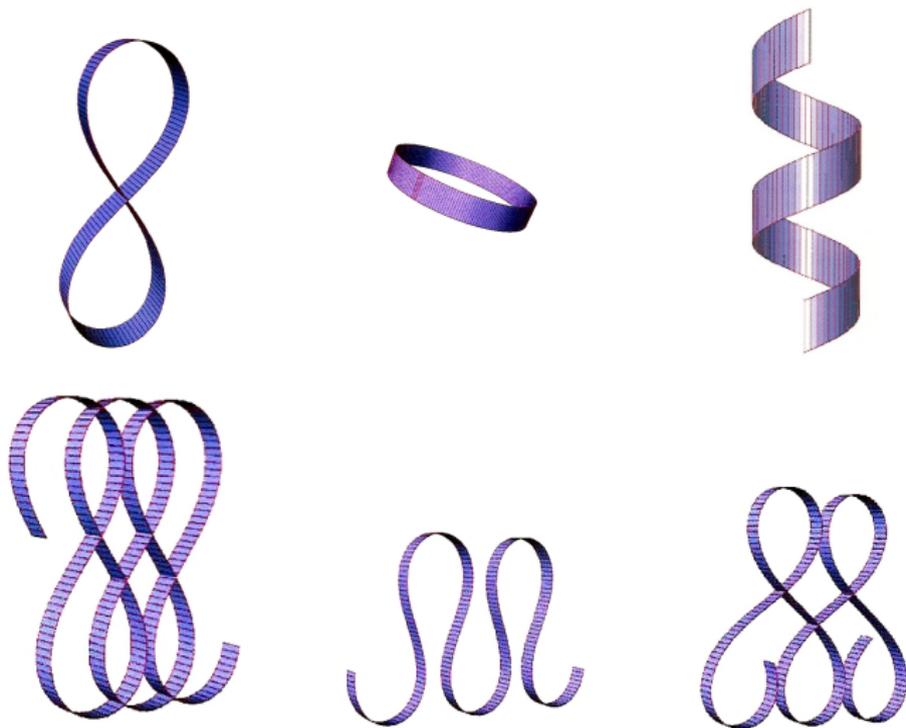


Computed in 2005 by Hagan, subsequently corrected by Rominger and Chubelaschwili

$$\begin{aligned}0 &= (\kappa'(1 + \lambda^2)^2 + 2\kappa(1 + \lambda^2)\lambda\lambda')' \\ &+ \frac{\kappa}{2}(\kappa^2(1 + \lambda^2)^2 - \mu) \\ &+ \lambda\kappa(\kappa^2(1 + \lambda^2)^2\lambda + (\frac{\kappa'}{\kappa}(1 + \lambda^2)2\lambda))' + (1 + \lambda^2)2\lambda)'' \\ 0 &= (\kappa^2(1 + \lambda^2)^2\lambda + (\frac{\kappa'}{\kappa}(1 + \lambda^2)2\lambda)' + ((1 + \lambda^2)2\lambda)'')' \\ &+ \kappa\lambda(\kappa'(1 + \lambda^2)^2 + 2\kappa(1 + \lambda^2)\lambda\lambda')\end{aligned}$$



Planar elastic curves and helices yield elastic strips with $\lambda = \text{const.}$





Other periodic examples can be found by numerical search (Rominger 2007).





Theorem (Chubelaschwili 2009): A strip is elastic \Leftrightarrow
the *force vector*

$$\begin{aligned} \mathbf{b} &= \frac{1}{2}(\kappa^2(1 + \lambda^2)^2 + \mu) \mathbf{T} \\ &+ (\kappa'(1 + \lambda^2)^2 + 2\kappa(1 + \lambda^2)\lambda\lambda') \mathbf{N} \\ &- (\kappa^2(1 + \lambda^2)^2\lambda + (\frac{\kappa'}{\kappa}(1 + \lambda^2)2\lambda)' + ((1 + \lambda^2)2\lambda)'') \mathbf{B} \end{aligned}$$

is constant.



Theorem (Chubelaschwili 2009): For an elastic strip
the *torque vector*

$$\begin{aligned} \mathbf{a} &= 2\kappa\lambda(1+\lambda^2)\mathbf{T} \\ &+ \frac{1}{\kappa}(2\kappa\lambda(1+\lambda^2))'\mathbf{N} \\ &+ \kappa(1+\lambda^2)(1-\lambda^2)\mathbf{B} \\ &- \mathbf{b} \times \boldsymbol{\gamma} \end{aligned}$$

is constant.



- ▶ The force \mathbf{b} and the torque \mathbf{a} have to be applied to the end point of the strip to keep it in equilibrium.
- ▶ They come from the boundary terms of the first variation formula used to derive the Euler-Lagrange equations.

Definition: An elastic strip is called *force-free* if the force vector \mathbf{b} vanishes.

- ▶ For force-free elastic strips the bending energy is critical even if the end point of γ is allowed to move, only the frame at the end point is held fixed.



- ▶ No condition on end point of γ
 \rightsquigarrow variational problem for the tangent image \mathbf{T}
- ▶ For a force-free elastic strip the Lagrange multiplier μ is positive
 \rightsquigarrow normalize to 1 by scaling.
- ▶ \rightsquigarrow We are looking for critical points of

$$\tilde{E} = 1/2 \int_0^L (\kappa^2(1 + \lambda^2)^2 + 1) ds.$$



- ▶ For a force-free elastic strip κ does not vanish \rightsquigarrow tangent image \mathbf{T} is a regular curve in S^2 of curvature λ :

$$\mathbf{T}' = \quad \quad + \kappa \mathbf{N}$$

$$\mathbf{N}' = -\kappa \mathbf{T} \quad \quad + \kappa \lambda \mathbf{B}$$

$$\mathbf{T}' = \quad \quad - \kappa \lambda \mathbf{N}$$

- ▶ $d\tilde{s} = \kappa ds \rightsquigarrow \gamma$ can be reconstructed from an arclength parametrization of \mathbf{T} as

$$\gamma(\tilde{s}) = \int_0^{\tilde{s}} \frac{1}{\kappa} \mathbf{T} d\tilde{s}.$$



Theorem: Let $\mathbf{T}: [0, \tilde{L}] \rightarrow S^2$ be an arclength parametrized curve with curvature λ . Then among all strips $\gamma: [0, \tilde{L}] \rightarrow \mathbb{R}^3$ with tangent image \mathbf{T} the one given by

$$\gamma(\tilde{s}) = \int_0^{\tilde{s}} (1 + \lambda^2) \mathbf{T} d\tilde{s}$$

has minimal Sadowski functional \tilde{E} and

$$\tilde{E} = \int_0^{\tilde{L}} (1 + \lambda^2) d\tilde{s}.$$

γ has curvature

$$\kappa = 1/(1 + \lambda)^2.$$



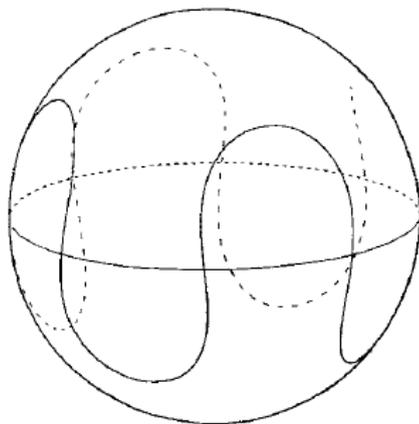
Corollary: The tangent images of force-free elastic strips are elastic curves in S^2 , in fact critical points of

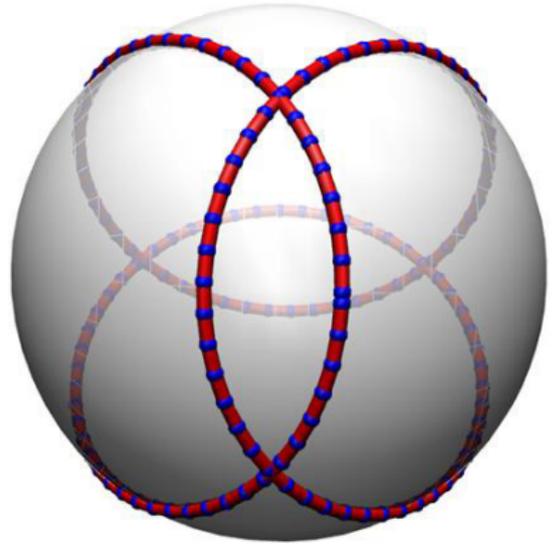
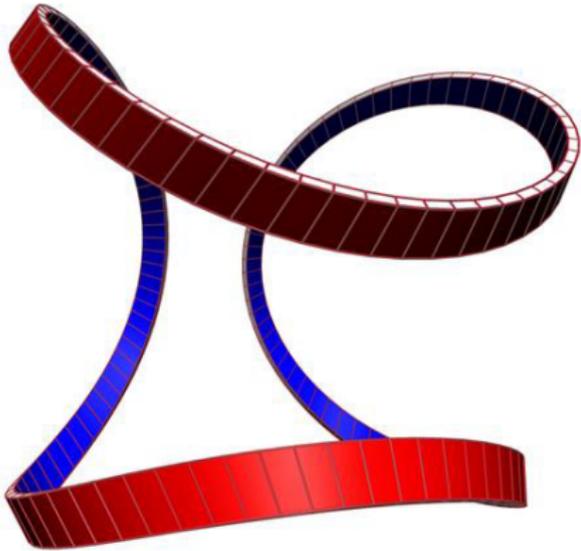
$$\int_0^{\tilde{L}} (1 + \lambda^2) d\tilde{s}.$$

Conversely, for any such spherical curve \mathbf{T} the space curve

$$\gamma = \int (1 + \lambda^2) \mathbf{T} d\tilde{s}$$

defines a force-free elastic strip.







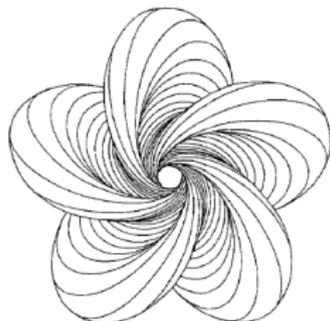
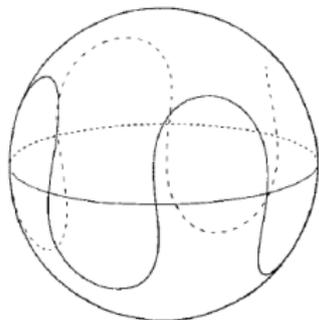
- ▶ All possible adapted frames (lifted to S^3) define the *frame cylinder*

$$F : [0, L] \times S^1 \rightarrow S^3$$

of a space curve γ .

- ▶ F is the preimage of the tangent image \mathbf{T} under the Hopf map $S^3 \rightarrow S^2$.

Corollary: The frame cylinder of a force-free elastic strip is Willmore in S^3 .





Given

- ▶ any arclength-parametrized spherical elastic curve $B : [0, L] \rightarrow S^2$ without inflection points
- ▶ its unit normal $T = B \times B'$
- ▶ its curvature λ .

Then

$$\gamma : [0, L] \rightarrow \mathbb{R}^3$$

$$\gamma = \int \left(1 + \frac{1}{\lambda^2}\right) T$$

defines an elastic strip.



Momentum strips are never closed.

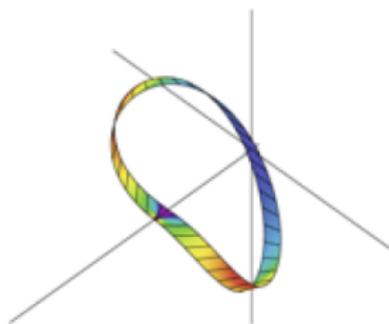


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At any point $\gamma(s)$ of a force-free elastic strip the following are equivalent:

- ▷ The rulings make an angle of 45° with γ .
- ▷ The curvature of the tangent image satisfies $\lambda(s) = 1$.
- ▷ A glueing construction is possible where the curvature κ is discontinuous but nevertheless we still have an elastic strip in the sense of balanced force \mathbf{b} and torque \mathbf{a} .





- ▶ Closed solutions by cut and paste?
- ▶ Stable ones?
- ▶ Moebius band?
- ▶ Other integrable classes of elastic strips?
- ▶ Maybe all elastic strips come from an integrable system?

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