

Polygonal smoke

Ulrich Pinkall jointly with Steffen Weißmann

DFG Research Center MATHEON Mathematics for key technologies



May 18, DDG 2010, Tennessee

▲□▶ ▲□▶ ▲ 三▶ ▲ 三 ● の < @



Smoke





Smoke rings





- Claim: The whole smoke can be modelled as a collection of entangled smoke rings.
- Smoke rings move on their own, but they also interact.
- Interaction can even imply a topology change (reconnection)







- Claim: The whole smoke can be modelled as a collection of entangled smoke rings.
- Smoke rings move on their own, but they also interact.
- Interaction can even imply a topology change (reconnection)







- Claim: The whole smoke can be modelled as a collection of entangled smoke rings.
- Smoke rings move on their own, but they also interact.
- Interaction can even imply a topology change (reconnection)







Vortex filaments



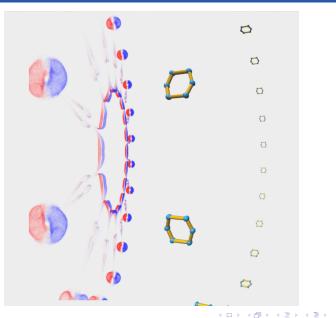






Polygonal Smoke





3



$\triangleright\,$ A velocity vector field v is uniquely determined by its vorticity

 $\omega = \operatorname{curl} \mathbf{v}$

 \triangleright v is given by the Biot-Savart formula:

$$v(x) = \frac{1}{4\pi} \int_{\mathbb{R}^3} \omega(y) \times \frac{x - y}{|x - y|^3} \, dy$$

 \triangleright In an ideal fluid ω flows with the velocity v it generates:

$$\dot{\omega} = \operatorname{curl}(\mathbf{v} \times \omega) = [\omega, \mathbf{v}]$$



\triangleright A velocity vector field v is uniquely determined by its vorticity

 $\omega = \operatorname{curl} \mathbf{v}$

 \triangleright v is given by the Biot-Savart formula:

$$u(x) = rac{1}{4\pi} \int_{\mathbb{R}^3} \omega(y) imes rac{x-y}{|x-y|^3} \, dy$$

 \triangleright In an ideal fluid ω flows with the velocity v it generates:

$$\dot{\omega} = \operatorname{curl}(\mathbf{v} \times \omega) = [\omega, \mathbf{v}]$$



\triangleright A velocity vector field v is uniquely determined by its vorticity

 $\omega = \operatorname{curl} \mathbf{v}$

 \triangleright v is given by the Biot-Savart formula:

$$v(x) = rac{1}{4\pi}\int_{\mathbb{R}^3}\omega(y) imesrac{x-y}{|x-y|^3}\,dy$$

 $\triangleright\,$ In an ideal fluid ω flows with the velocity v it generates:

$$\dot{\omega} = \operatorname{curl}(\mathbf{v} \times \omega) = [\omega, \mathbf{v}]$$

< 3 > < 3 >

Origin of vorticity



Away from obstacles vorticity is neither created nor destroyed

▷ Just swept along with the flow

- All vorticity originates at the boundaries of obstacles
- Kaffeelöffelexperiment by Felix Klein



Origin of vorticity



- Away from obstacles vorticity is neither created nor destroyed
- $\,\triangleright\,$ Just swept along with the flow
- All vorticity originates at the boundaries of obstacles
- Kaffeelöffelexperiment by Felix Klein

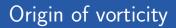


Origin of vorticity



- Away from obstacles vorticity is neither created nor destroyed
- $\,\triangleright\,$ Just swept along with the flow
- All vorticity originates at the boundaries of obstacles
- Kaffeelöffelexperiment by Felix Klein



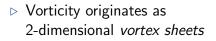




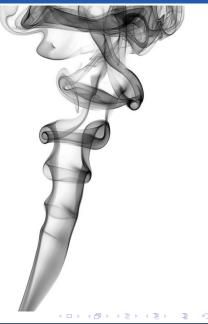
- Away from obstacles vorticity is neither created nor destroyed
- $\,\triangleright\,$ Just swept along with the flow
- All vorticity originates at the boundaries of obstacles
- Kaffeelöffelexperiment by Felix Klein







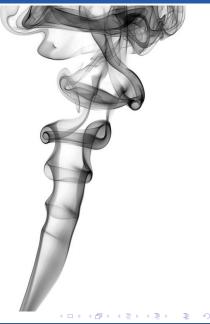
 Vortex sheets roll up into 1-dimensional structures ("smoke rings")





Vorticity originates as
 2-dimensional vortex sheets

 Vortex sheets roll up into 1-dimensional structures ("smoke rings")



- Airplane rides on a giant vortex ring
- Extends
 back to
 where it
 took off
- Vorticity concentrated on a filament



- Airplane rides on a giant vortex ring
- Extends back to where it took off
- Vorticity concentrated on a filament



- Airplane
 rides on a
 giant
 vortex ring
- Extends back to where it took off
- Vorticity concentrated on a filament





 Suppose all vorticity is concentrated in a small tube of radius *R* around a space curve γ

(like water flowing through the tube).

 \triangleright Then away from γ the velocity field is given by

$$v(x) = K \oint \frac{\gamma' \times (x - \gamma)}{|x - \gamma|^3}$$





▷ Suppose all vorticity is concentrated in a small tube of radius *R* around a space curve *γ*

(like water flowing through the tube).

 $\triangleright\,$ Then away from γ the velocity field is given by

$$\mathbf{v}(\mathbf{x}) = \mathbf{K} \oint rac{\gamma' imes (\mathbf{x} - \gamma)}{|\mathbf{x} - \gamma|^3}$$





$\triangleright\,$ Evolution of $\gamma:$ Evaluate velocity v on $\gamma \rightsquigarrow$

 $\dot{\gamma} pprox \mathit{C_f} \ \mathit{K} \ \log(\mathit{R}) \ \gamma' imes \gamma''$

 $\triangleright \text{ Scale down } K \text{ as } R \rightarrow 0 \rightsquigarrow \text{ smoke ring flow}$

⊳ da Rios and Levi-Civita 1906

 Integrable system equivalent to the non-linear Schroedinger equation (Hashimoto 1972)





 $\triangleright\,$ Evolution of $\gamma:$ Evaluate velocity v on $\gamma \rightsquigarrow$

 $\dot{\gamma} pprox \mathcal{C}_{f} \ \mathcal{K} \ \log(\mathcal{R}) \ \gamma' imes \gamma''$

\triangleright Scale down K as $R \rightarrow 0 \rightsquigarrow$ smoke ring flow

▷ da Rios and Levi-Civita 1906

 Integrable system equivalent to the non-linear Schroedinger equation (Hashimoto 1972)





 $\triangleright\,$ Evolution of $\gamma:$ Evaluate velocity v on $\gamma \rightsquigarrow$

 $\dot{\gamma} pprox \mathcal{C}_{f} \ \mathcal{K} \ \log(\mathcal{R}) \ \gamma' imes \gamma''$

- \triangleright Scale down *K* as $R \rightarrow 0 \rightsquigarrow$ smoke ring flow
- b da Rios and Levi-Civita 1906
- Integrable system equivalent to the non-linear Schroedinger equation (Hashimoto 1972)





 $\triangleright\,$ Evolution of $\gamma:$ Evaluate velocity v on $\gamma \rightsquigarrow$

 $\dot{\gamma} \approx C_f \ K \ \log(R) \ \gamma' \times \gamma''$

 $\triangleright\,$ Scale down K as $R\to 0\rightsquigarrow\,$ smoke ring flow

- b da Rios and Levi-Civita 1906
- Integrable system equivalent to the non-linear Schroedinger equation (Hashimoto 1972)





 $\dot{\gamma} = \gamma' \times \gamma''$ \triangleright

- Curve moves orthogonal to its osculating plane.
- Speed is proportional to the curvature.
- ▷ Length is constant.
 - Area vector is constant.





$$\triangleright \quad \dot{\gamma} = \gamma' \times \gamma''$$

- Curve moves orthogonal to its osculating plane.
- Speed is proportional to the curvature.
- ▷ Length is constant.
 - Area vector is constant.





$$\triangleright \quad \dot{\gamma} = \gamma' \times \gamma''$$

- Curve moves orthogonal to its osculating plane.
- Speed is proportional to the curvature.
- ▷ Length is constant.





$$\triangleright \quad \dot{\gamma} = \gamma' \times \gamma''$$

- Curve moves orthogonal to its osculating plane.
- Speed is proportional to the curvature.
- ▷ Length is constant.





$$\triangleright \quad \dot{\gamma} = \gamma' \times \gamma''$$

- Curve moves orthogonal to its osculating plane.
- Speed is proportional to the curvature.
- ▷ Length is constant.





$$\triangleright \quad \dot{\gamma} = \gamma' \times \gamma''$$

- Curve moves orthogonal to its osculating plane.
- Speed is proportional to the curvature.
- ▷ Length is constant.



. . .



Area vector of a polygon

▷ For a space curve:

$$A = \frac{1}{2} \oint \gamma \times \gamma'$$

▷ For a closed polygon:

$$A = \frac{1}{2} \sum_{i=1}^{n} \gamma_i \times \gamma_{i+1}$$

▷ b a unit vector → ⟨A, b⟩ is the algebraic area of the orthogonal projection of γ onto a plane with normal vector b





Area vector of a polygon

▷ For a space curve:

$$A = \frac{1}{2} \oint \gamma \times \gamma'$$

▷ For a closed polygon:

$$\mathsf{A} = \frac{1}{2}\sum_{i=1}^n \gamma_i \times \gamma_{i+1}$$

▷ b a unit vector → ⟨A, b⟩ is the algebraic area of the orthogonal projection of γ onto a plane with normal vector b





Area vector of a polygon

▷ For a space curve:

$$A = \frac{1}{2} \oint \gamma \times \gamma'$$

▷ For a closed polygon:

$$A = \frac{1}{2} \sum_{i=1}^{n} \gamma_i \times \gamma_{i+1}$$

b a unit vector → ⟨A, b⟩ is the algebraic area of the orthogonal projection of γ onto a plane with normal vector b





Area vector of a polygon

▷ For a space curve:

$$A = \frac{1}{2} \oint \gamma \times \gamma'$$

▷ For a closed polygon:

$$A = \frac{1}{2} \sum_{i=1}^{n} \gamma_i \times \gamma_{i+1}$$

b a unit vector → ⟨A, b⟩ is the algebraic area of the orthogonal projection of γ onto a plane with normal vector b



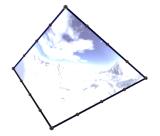


▷ A quadrilateral $\gamma_0, \gamma_1, \gamma_2, \gamma_3$ in \mathbb{R}^3 is called a *skew parallelogram* of twist τ if the difference vector

$$V = \frac{\gamma_3 + \gamma_1}{2} - \frac{\gamma_2 + \gamma_0}{2}$$

between the centers of its diagonals is a multiple of its area vector:

$$V = \tau A$$



 Opposite sides of a skew parallelogram have the same length.

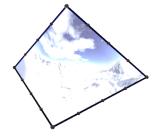


A quadrilateral γ₀, γ₁, γ₂, γ₃ in ℝ³ is called a *skew parallelogram* of twist τ if the difference vector

$$V = \frac{\gamma_3 + \gamma_1}{2} - \frac{\gamma_2 + \gamma_0}{2}$$

between the centers of its diagonals is a multiple of its area vector:

$$V = \tau A$$



 Opposite sides of a skew parallelogram have the same length.



A quadrilateral γ₀, γ₁, γ₂, γ₃ in ℝ³ is called a *skew parallelogram* of twist τ if the difference vector

$$V = \frac{\gamma_3 + \gamma_1}{2} - \frac{\gamma_2 + \gamma_0}{2}$$

between the centers of its diagonals is a multiple of its area vector:

$$V = \tau A$$



Opposite sides of a skew parallelogram have the same length.



 A polygon η is called a Darboux transform with rod-length ρ and twist τ if all quadrilaterals

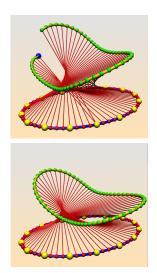
$$\eta_i \quad \eta_{i+1}$$

 $\gamma_i \quad \gamma_{i+1}$
are skew parallelograms with

$$|\eta_i - \gamma_i| = \rho$$
 and twist τ .

▷ For generic ρ, τ every closed polygon has exactly two closed Darboux transforms.

⊳ Hoffmann 2000





 A polygon η is called a Darboux transform with rod-length ρ and twist τ if all quadrilaterals

$$\eta_i \quad \eta_{i+1}$$

 $\gamma_i \quad \gamma_{i+1}$
are skew parallelograms with

$$|\eta_i - \gamma_i| = \rho$$
 and twist τ .

 $\triangleright\,$ For generic ρ,τ every closed polygon has exactly two closed Darboux transforms.



 A polygon η is called a Darboux transform with rod-length ρ and twist τ if all quadrilaterals

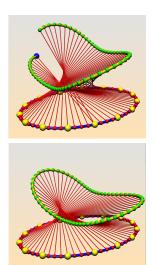
$$\eta_i \quad \eta_{i+1}$$

 $\gamma_i \quad \gamma_{i+1}$
are skew parallelograms with

$$|\eta_i - \gamma_i| = \rho$$
 and twist τ .

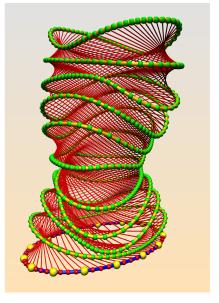
 $\triangleright\,$ For generic ρ,τ every closed polygon has exactly two closed Darboux transforms.

Hoffmann 2000



Closed Darboux transforms of a closed polygon have the same:

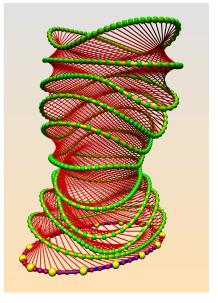
length Area vecto



Closed Darboux transforms of a closed polygon have the same:

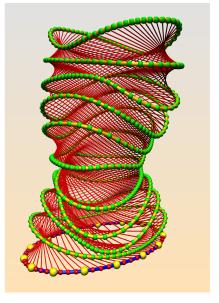
length

Area vector



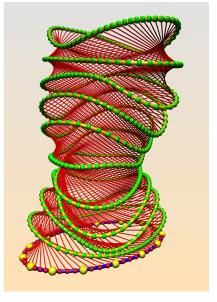
Closed Darboux transforms of a closed polygon have the same:

length Area vector



Closed Darboux transforms of a closed polygon have the same:

length Area vector

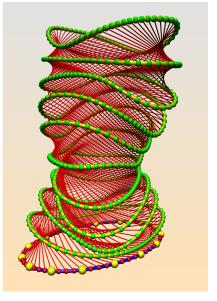




Closed Darboux transforms of a closed polygon have the same:

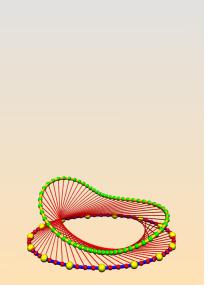
length

Area vector



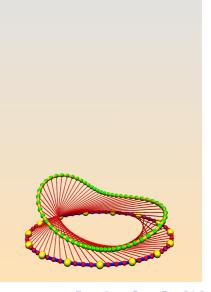


- \triangleright Using twists τ and $-\tau$ in an alternating fashion preserves reflectional symmetries.
- ▷ Forget the odd iterations.
- Excellent discrete version of the smoke-ring flow.



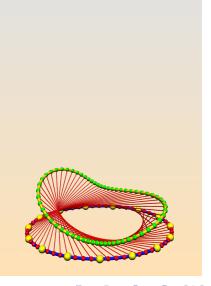


- \triangleright Using twists τ and $-\tau$ in an alternating fashion preserves reflectional symmetries.
- ▷ Forget the odd iterations.
- Excellent discrete version of the smoke-ring flow.



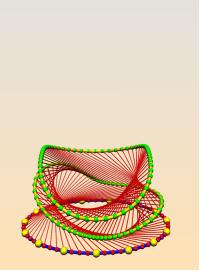


- \triangleright Using twists τ and $-\tau$ in an alternating fashion preserves reflectional symmetries.
- ▷ Forget the odd iterations.
- Excellent discrete version of the smoke-ring flow.



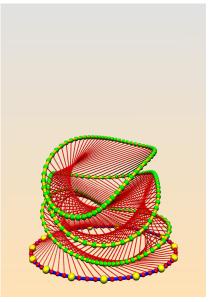


- \triangleright Using twists τ and $-\tau$ in an alternating fashion preserves reflectional symmetries.
- ▷ Forget the odd iterations.
- Excellent discrete version of the smoke-ring flow.



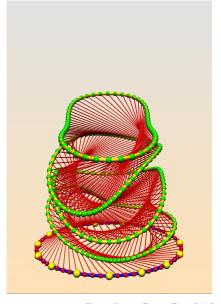


- \triangleright Using twists τ and $-\tau$ in an alternating fashion preserves reflectional symmetries.
- ▷ Forget the odd iterations.
- Excellent discrete version of the smoke-ring flow.



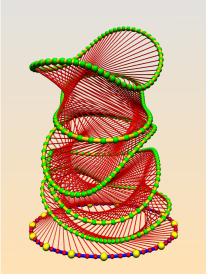


- \triangleright Using twists τ and $-\tau$ in an alternating fashion preserves reflectional symmetries.
- ▷ Forget the odd iterations.
- Excellent discrete version of the smoke-ring flow.



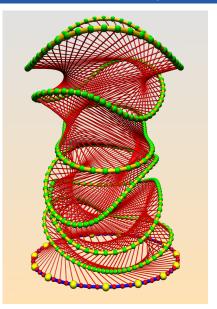


- \triangleright Using twists τ and $-\tau$ in an alternating fashion preserves reflectional symmetries.
- ▷ Forget the odd iterations.
- Excellent discrete version of the smoke-ring flow.





- \triangleright Using twists τ and $-\tau$ in an alternating fashion preserves reflectional symmetries.
- ▷ Forget the odd iterations.
- Excellent discrete version of the smoke-ring flow.





To make this practical, several features have to be included:

Interaction between thick vortex rings

Obstacles

 Vorticity generation at obstacle boundaries ("vortex shedding")

▷ Topology changes ("vortex reconnection")









To make this practical, several features have to be included:

- Interaction between thick vortex rings
- Obstacles
- Vorticity generation at obstacle boundaries ("vortex shedding")
- ▷ Topology changes ("vortex reconnection")









To make this practical, several features have to be included:

- Interaction between thick vortex rings
- \triangleright Obstacles
- Vorticity generation at obstacle boundaries ("vortex shedding")
- ▷ Topology changes ("vortex reconnection")









To make this practical, several features have to be included:

- Interaction between thick vortex rings
- Obstacles
- Vorticity generation at obstacle boundaries ("vortex shedding")

Description Topology changes ("vortex reconnection")









To make this practical, several features have to be included:

- Interaction between thick vortex rings
- Obstacles
- Vorticity generation at obstacle boundaries ("vortex shedding")
- ▷ Topology changes ("vortex reconnection")









See Steffen's talk on Thursday!



2